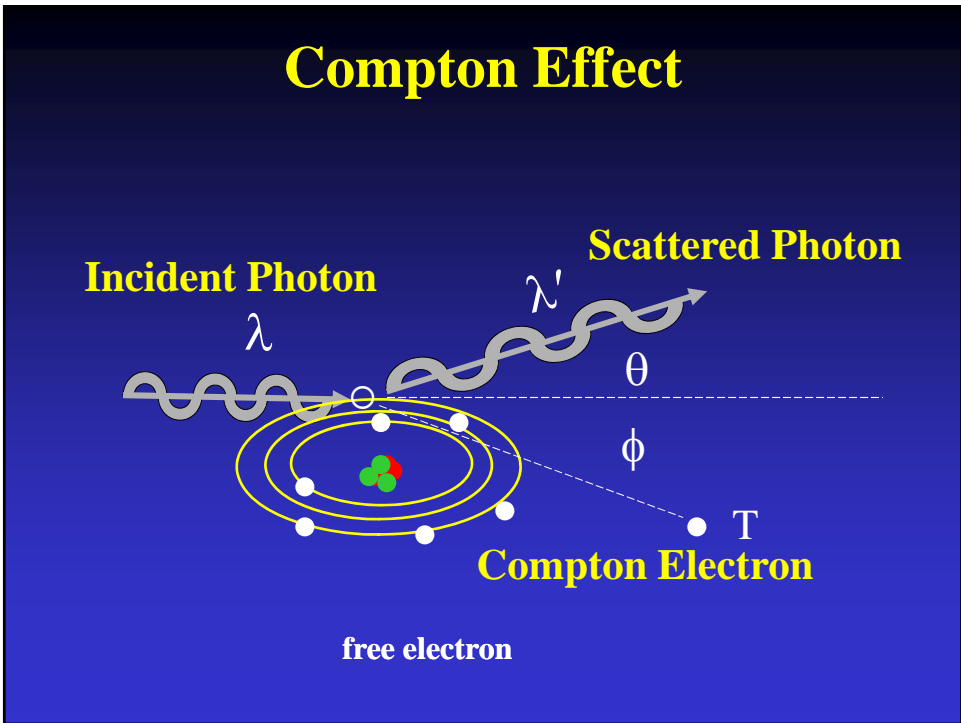
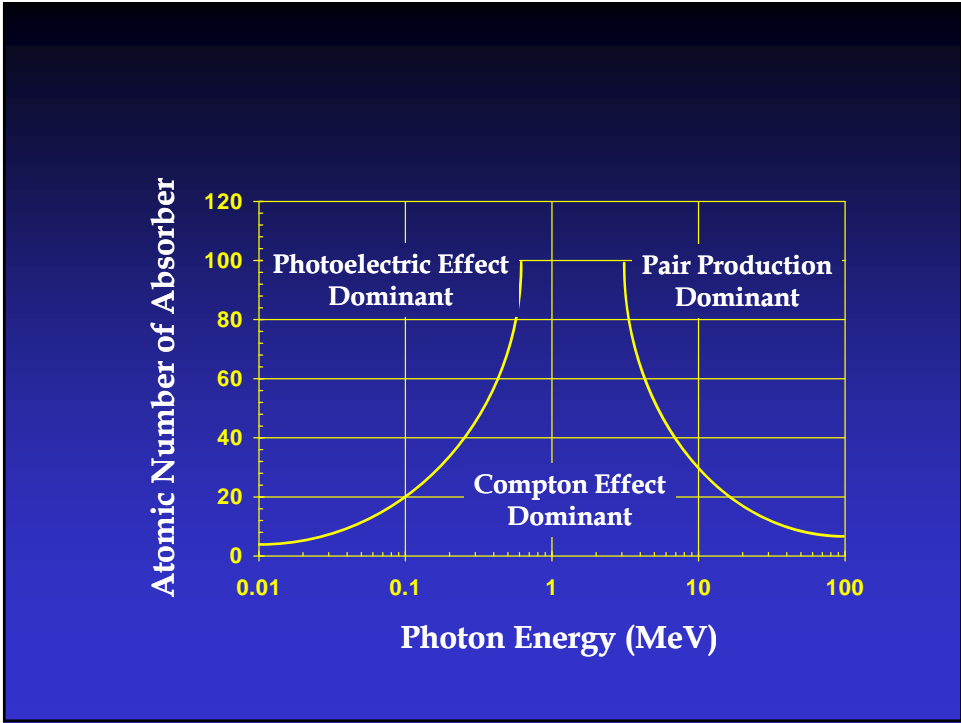


## **Interaction of Gamma- and X-ray in Matter**

### **Five Interactions**

- **Compton effect**
- **Photoelectric effect**
- **Pair production**
- **Rayleigh (coherent) scattering**
- **Photonuclear interactions**



## Kinematics

These equations are results of applying conservation of energy and momentum to the Compton interaction

$$h\nu' = \frac{h\nu}{1 + (h\nu/m_0c^2)(1 - \cos \varphi)}$$

$$T = h\nu - h\nu'$$

$$\cot \theta = \left(1 + \frac{h\nu}{m_0c^2}\right) \tan \left(\frac{\varphi}{2}\right)$$

$$h\nu' = \frac{h\nu}{1 + (h\nu/m_0c^2)(1 - \cos \varphi)}$$

$$T = h\nu - h\nu'$$

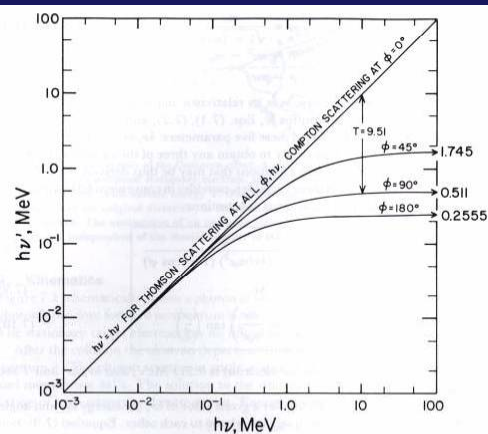


FIGURE 7.3. Graphical representation of the kinematic relationship of  $h\nu$ ,  $h\nu'$ , and  $T$  in the Compton effect, as described by Eqs. (7.8) and (7.9). Curves are shown only for  $\varphi = 0^\circ, 45^\circ, 90^\circ$  and  $180^\circ$ . Note that  $T$  is to be interpreted as the vertical separation of any  $\varphi$ -curve from the  $\varphi = 0^\circ$  diagonal. In the case shown ( $h\nu = 10$  MeV,  $\varphi = 90^\circ$ ),  $T = 9.51$  MeV.

$$\cot \theta = \left( 1 + \frac{h\nu}{m_0c^2} \right) \tan \left( \frac{\varphi}{2} \right)$$

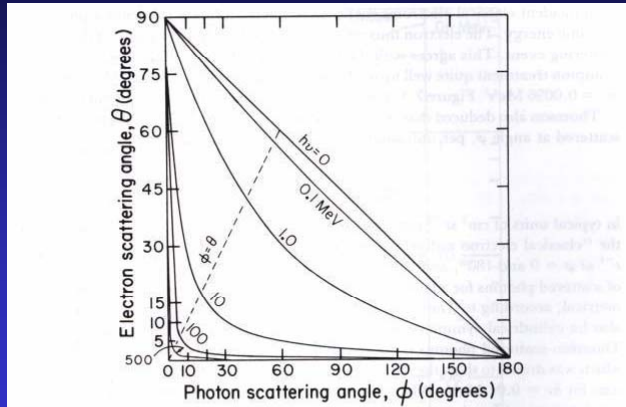


FIGURE 7.4. Relationship of the electron scattering angle  $\theta$  to the photon scattering angle  $\varphi$  in the Compton effect, from Eq. (7.10). Curves are shown for the incident photon energies 0, 0.1, 1.0, 10, 100, and 500 MeV. The dashed line is the locus where  $\theta = \varphi$ , when the electron and photon are scattered at equal angles on opposite sides of the incident photon's direction.

## Interaction cross section

Thomson Scattering

classical electron radius

$$\frac{d_e\sigma_0}{d\Omega_\varphi} = \frac{r_0^2}{2} (1 + \cos^2 \varphi)$$

Valid only for non-relativistic approximation

Klein-Nishina cross sections for photon scattering at angle  $\varphi$

$$\frac{d_e\sigma}{d\Omega_\varphi} = \frac{r_0^2}{2} \left( \frac{h\nu'}{h\nu} \right)^2 \left( \frac{h\nu}{h\nu'} + \frac{h\nu'}{h\nu} - \sin^2 \varphi \right)$$

Use of Dirac's relativistic theory of the electron

$h\nu = h\nu'$

$$\frac{d_e\sigma}{d\Omega_\phi} = \frac{r_0^2}{2} \left(\frac{h\nu'}{h\nu}\right)^2 \left(\frac{h\nu}{h\nu'} + \frac{h\nu'}{h\nu} - \sin^2 \phi\right)$$

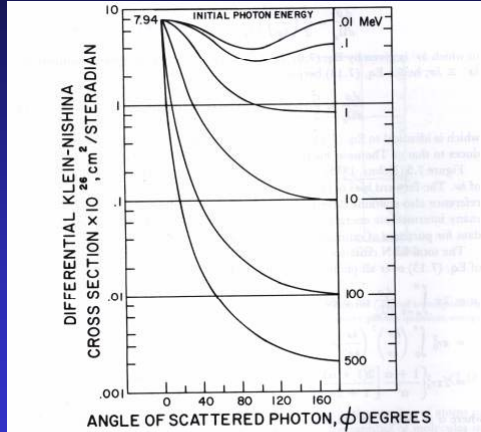


FIGURE 7.5. Differential Klein-Nishina cross section,  $d_e\sigma/d\Omega_\phi$  [see Eq. (7.13)] vs. angle  $\phi$  of the scattered photon, for  $h\nu = 0.01, 0.1, 1.0, 10, 100,$  and  $500$  MeV. This shows the angular distribution, per unit solid angle, of the scattered photons resulting from the Compton effect. (After Nelms, 1953.)

$$\begin{aligned} \sigma &= 2\pi \int_0^\pi \frac{d_e\sigma}{d\Omega_\phi} \sin \phi \, d\phi \\ &= \pi r_0^2 \int_0^\pi \left(\frac{h\nu'}{h\nu}\right)^2 \left(\frac{h\nu}{h\nu'} + \frac{h\nu'}{h\nu} - \sin^2 \phi\right) \sin \phi \, d\phi \\ &= 2\pi r_0^2 \left\{ \frac{1 + \alpha}{\alpha^2} \left[ \frac{2(1 + \alpha)}{1 + 2\alpha} - \frac{\ln(1 + 2\alpha)}{\alpha} \right] + \frac{\ln(1 + 2\alpha)}{2\alpha} - \frac{1 + 3\alpha}{(1 + 2\alpha)^2} \right\} \end{aligned}$$

$$\alpha = \frac{h\nu}{m_0c^2}$$

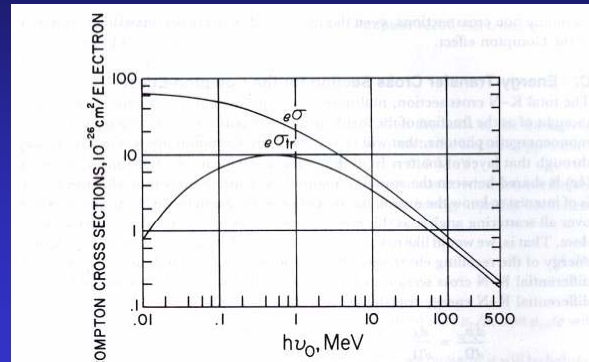


FIGURE 7.6. Klein-Nishina (Compton-effect) cross section per electron ( $\sigma$ ) and corresponding energy-transfer cross section per electron ( $\sigma_{tr}$ ) as a function of primary photon quantum energy  $h\nu$ . (After Nelms, 1953.)

$${}_e\sigma \propto Z^0$$

$${}_a\sigma = Z \cdot {}_e\sigma \quad (\text{cm}^2/\text{atom})$$

Avogadro's number

$$\frac{\sigma}{\rho} = \frac{N_A Z}{A} {}_e\sigma \quad (\text{cm}^2/\text{g})$$

Number of grams per mole of material

## Energy transfer cross section

Total K-N cross section, multiplied by a unit thickness of 1 e/cm<sup>2</sup> may be thought as the fraction of the incident energy fluence, carried by a beam of many monoenergetic photons

What is the overall fraction of  $h\nu$  that is given to the electrons, averaged over all scattering angles? This energy contributes to the kerma and therefore to dose.

$h\nu - h\nu'$

$$\frac{d_e\sigma_{\text{tr}}}{d\Omega_\varphi} = \frac{d_e\sigma}{d\Omega_\varphi} \cdot \frac{T}{h\nu}$$

$$\frac{\overline{T}}{h\nu} = \frac{e\sigma_{tr}}{e\sigma}$$

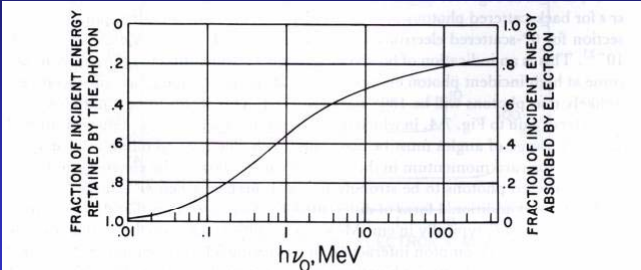


FIGURE 7.7. Mean fraction ( $\overline{T}/h\nu$ ) of the incident photon's energy given to the recoiling electron in Compton interactions, averaged over all angles (right ordinate). Also, mean fraction ( $h\nu'/h\nu$ ) of energy retained by the scattered photon (left ordinate).

Differential K-N cross section for electron scattering at angle  $\theta$ , per unit solid angle per electron

$$\frac{d_e\sigma}{d\Omega_\theta} = \frac{d_e\sigma}{d\Omega_\varphi} \cdot \frac{(1 + \alpha)^2 (1 - \cos \varphi)^2}{\cos^3 \theta}$$

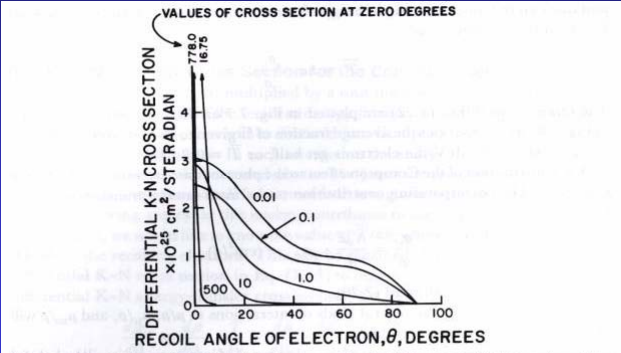


FIGURE 7.8. Differential Klein-Nishina cross section  $d_e\sigma/d\Omega_\theta$ , vs angle  $\theta$  of the scattered electron for  $h\nu = 0.01, 0.1, 1, 10,$  and  $500$  MeV. This shows the angular distribution, per unit solid angle, of the recoil electrons resulting from the Compton effect. (After Nelms, 1953.)

## Energy distribution of electrons, averaged over all scattering angles

$$\frac{d\sigma}{dT} = \frac{\pi r_0^2 m_0 c^2}{(h\nu')^2} \times \left\{ \left[ \frac{m_0 c^2 T}{(h\nu)^2} \right]^2 + 2 \left[ \frac{h\nu'}{h\nu} \right]^2 + \frac{h\nu'}{(h\nu)^3} [(T - m_0 c^2)^2 - (m_0 c^2)^2] \right\}$$

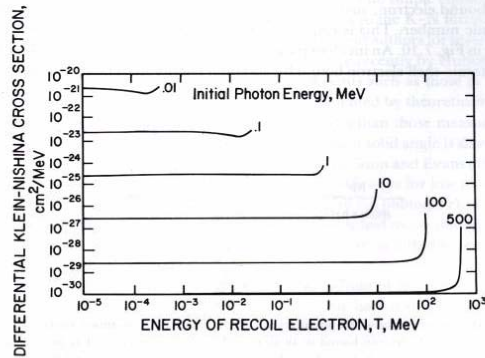
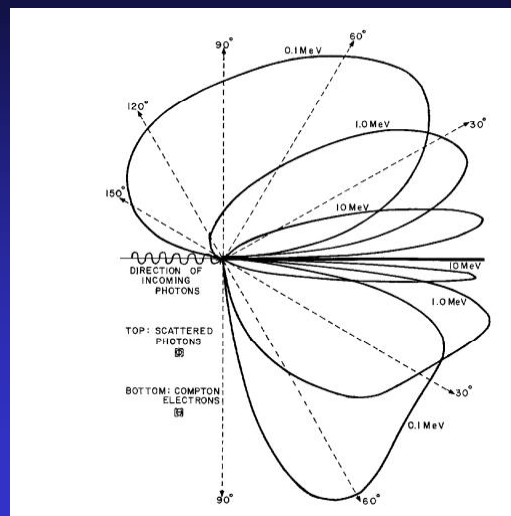


FIGURE 7.9. Differential Klein-Nishina cross section  $d\sigma/dT$  expressing the initial energy spectrum of Compton recoiling electrons. (After Nelms, 1953.)

## Compton Effect

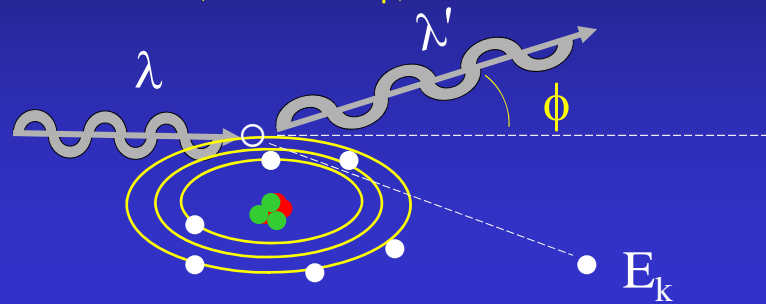


## Compton Effect

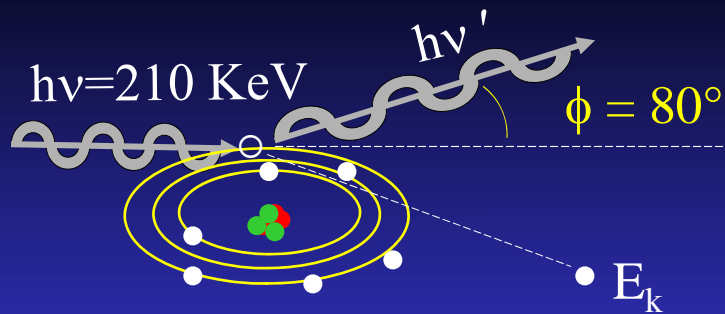
$$h\nu(\text{KeV}) = 12.4 / \lambda(\text{\AA})$$

$$\lambda' = \lambda + \Delta\lambda$$

$$\Delta\lambda = 0.0243(1 - \cos \phi)$$



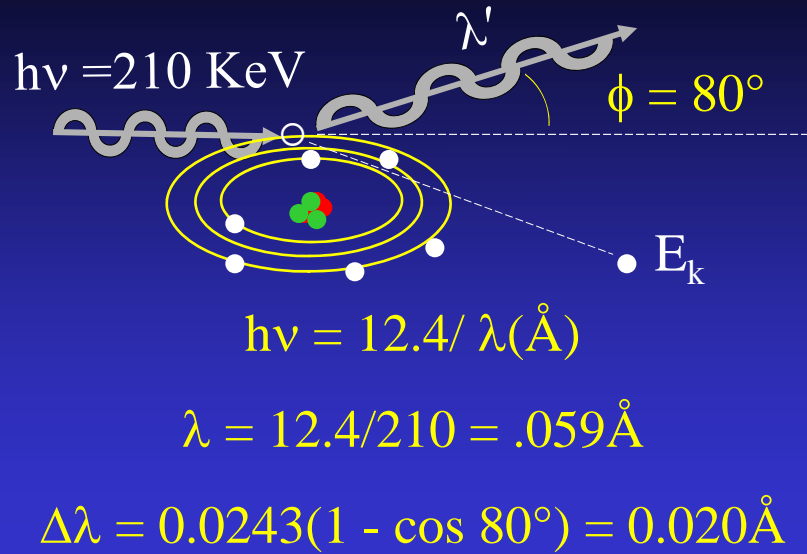
## Compton Effect



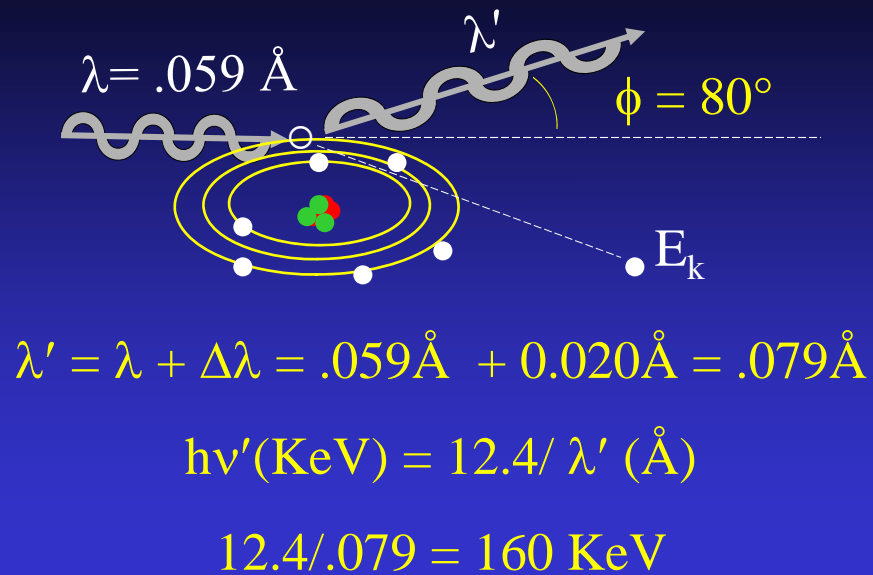
What is  $h\nu'$  ?

What is  $E_k$  ?

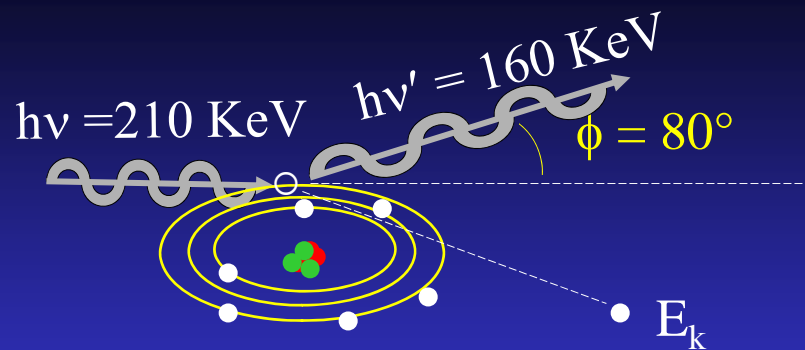
## Compton Effect



## Compton Effect



## Compton Effect



$$E_k = h\nu - h\nu'$$

$$E_k = 210 \text{ KeV} - 160 \text{ KeV}$$

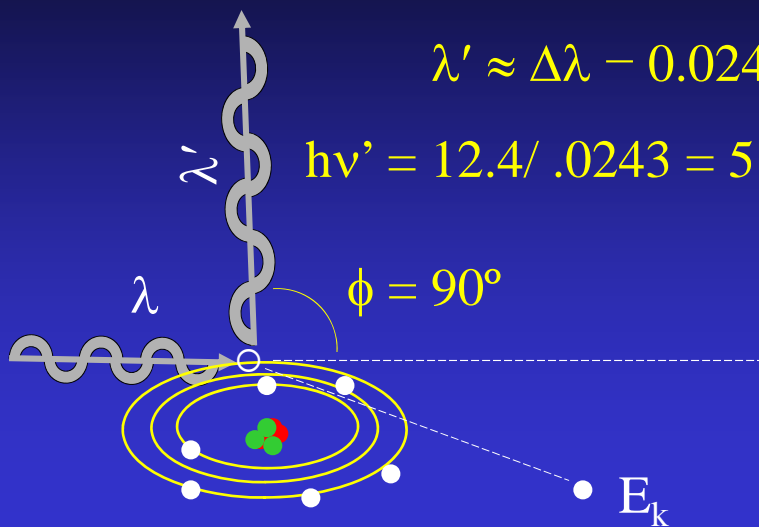
$$E_k = 50 \text{ KeV}$$

## Compton Effect

$$\cos 90^\circ = 0$$

$$\lambda' \approx \Delta\lambda - 0.0243$$

$$h\nu' = 12.4 / .0243 = 511 \text{ KeV}$$

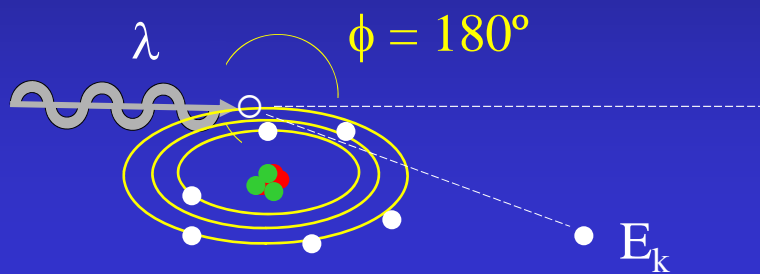


## Compton Effect

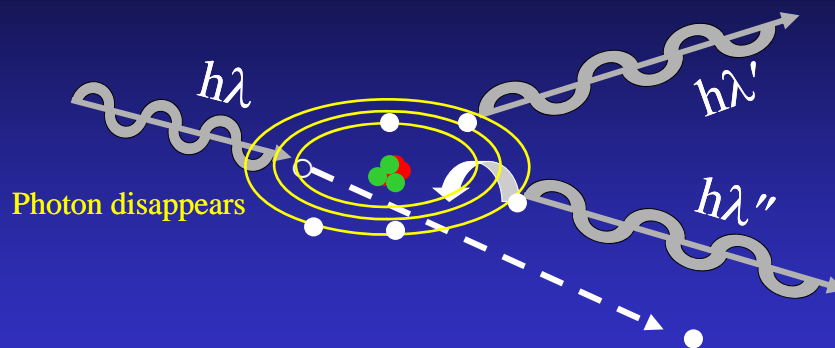
$$\cos 180^\circ = -1$$

$$\lambda' \approx \Delta\lambda - 0.0243 \cdot 2$$

$$h\nu' = 12.4 / .0486 = 255 \text{ KeV}$$

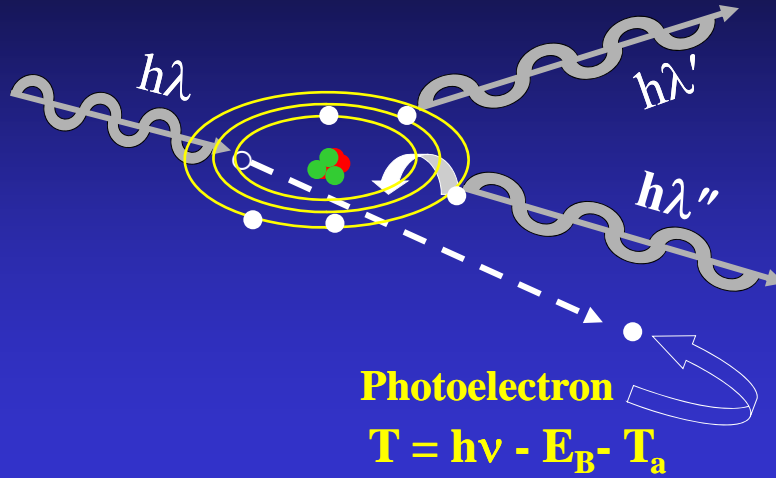


## Photoelectric Effect

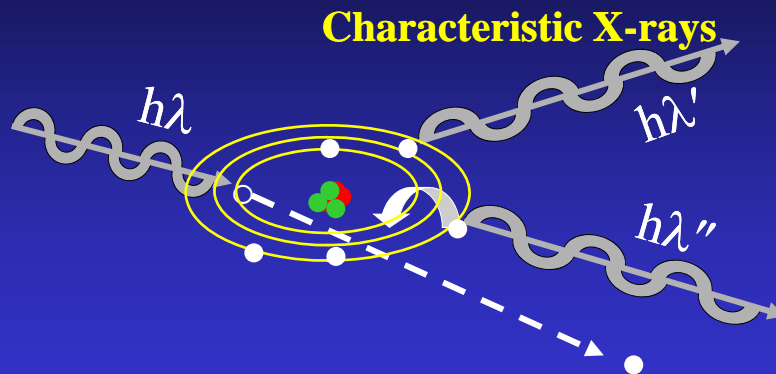


**X-ray photon interacts with an atom and ejects one of orbital electrons.**

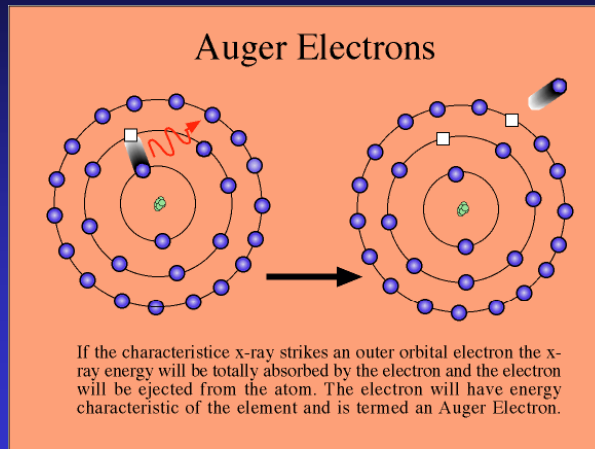
## Photoelectric Effect



## Photoelectric Effect

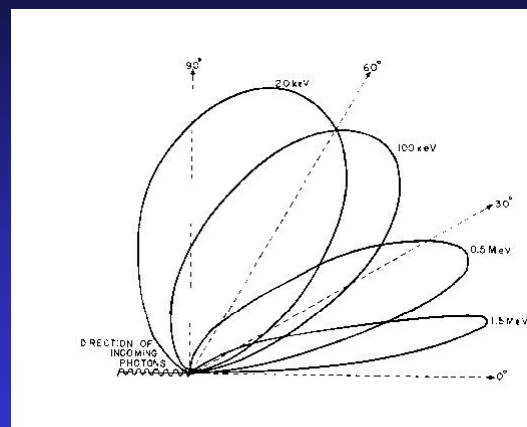


## Photoelectric Effect



Auger electrons are emitted as a competing process for characteristic x-rays

## Photoelectrons



## Interaction cross section

There are no simple formulas for photoelectric cross sections, this is because of complications due to the binding energy of the electron. Tables are based on experimental results and supplemented by theoretically assisted interpolations for other energies and absorbing media.

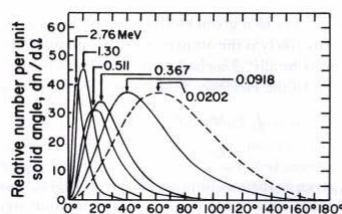


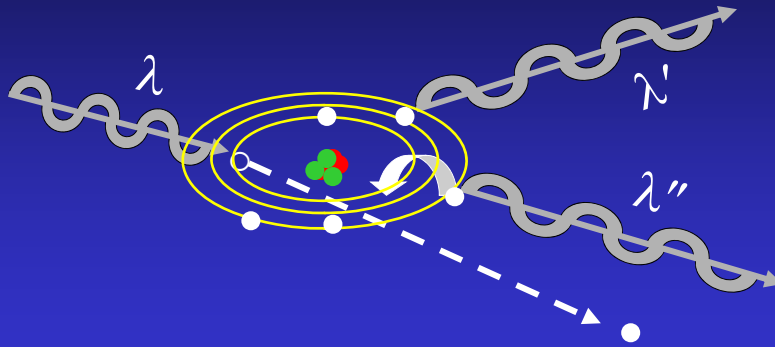
FIGURE 7.11. Directional distribution of photoelectrons per unit solid angle, for energies as labeled on the curves. The curve areas are not normalized to each other. [After Davison and Evans (1952). Reproduced with permission of R. D. Evans and the American Physical Society.]

$${}_a\tau \cong k \frac{Z^n}{(h\nu)^m} \quad (\text{cm}^2/\text{atom})$$

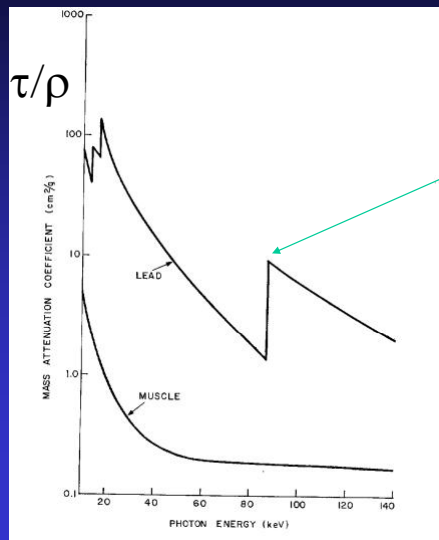
n = 4 at 0.1 MeV → 4.6 at 3 MeV  
m = 3 at 0.1 MeV → 1 at 5 MeV

$$\frac{\tau}{\rho} \propto \left( \frac{Z}{h\nu} \right)^3 \quad (\text{cm}^2/\text{g})$$

# Photoelectric Effect

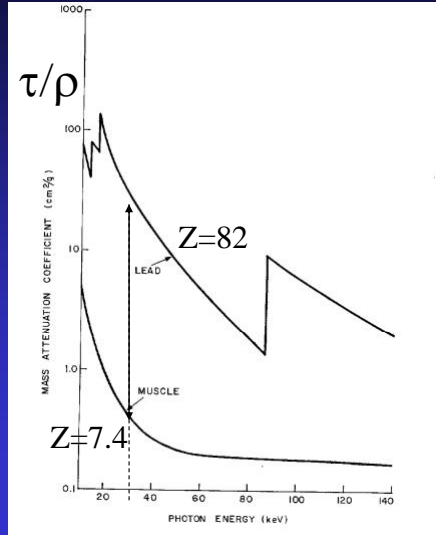


# Photoelectric Effect



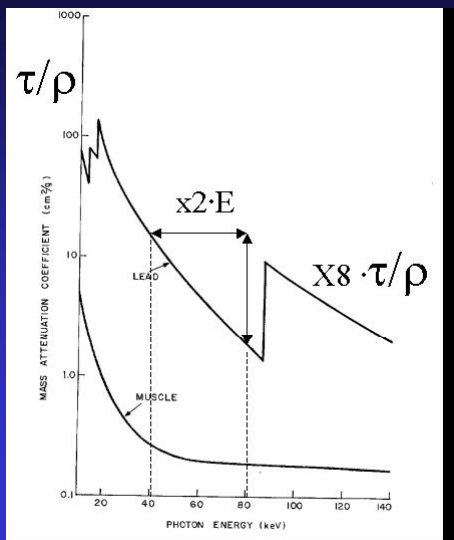
Characteristic X-rays

# Photoelectric Effect



$$\tau/\rho \propto Z^3/E^3$$

# Photoelectric Effect



$$\tau/\rho \propto Z^3/E^3$$

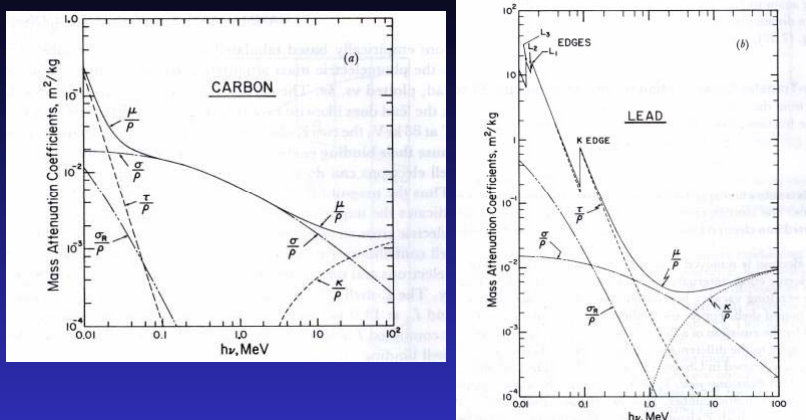


FIGURE 7.13. Mass attenuation coefficients for carbon (a) and lead (b).  $\tau/\rho$  indicates the contribution of the photoelectric effect,  $\sigma/\rho$  is that of the Compton effect,  $\kappa/\rho$  that of pair production, and  $\sigma_R/\rho$  that of Rayleigh (coherent) scattering.  $\mu/\rho$  is their sum, which is closely approximated in Pb by the  $\tau/\rho$  curve below  $h\nu = 0.1$  MeV (From data of Hubbell, 1969).

## Energy-transfer cross section

$$\frac{T}{h\nu} = \frac{h\nu - E_b}{h\nu}$$

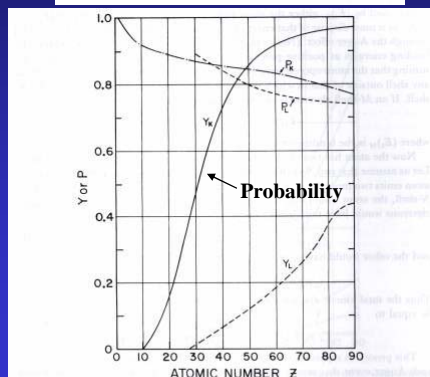


FIGURE 7.14. Fluorescence yield ( $Y_K$ ) and fractional participation in the photoelectric effect ( $P_K$ ) by K- and L-shell electrons (see text).  $P_K$  and  $P_L$  was calculated from tables of Hubbell (1969) and McMaster et al. (1969);  $Y_K$  from Lederer and Shirley (1979); and  $Y_L$  from Burhop (1952).

$P_K$  is the fraction of all photoelectric interactions that occur in the K shell.

$P_K Y_K$  is the fraction of all photoelectric events in which a K fluorescence x-ray is emitted by the atom.

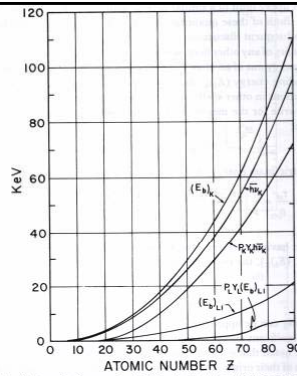


FIGURE 7.15. Electron binding energies ( $E_{K\alpha}$  in the  $K$ -shell and  $(E_{L\alpha})$  in the  $L$ -shell; weighted mean fluorescence x-ray energy  $h\bar{\nu}_K$  in the  $K$ -shell; and the products  $P_K Y_K h\bar{\nu}_K$  and  $P_L Y_L h\bar{\nu}_L$ . The latter provides an upper-limit estimate of  $P_L Y_L h\bar{\nu}_L$ . Taken or derived from tables by Lederer and Shirley (1979).

$$\frac{\tau_{tr}}{\rho} = \frac{\tau}{\rho} \left[ \frac{h\nu - P_K Y_K h\bar{\nu}_K - (1 - P_K) P_L Y_L h\bar{\nu}_L}{h\nu} \right]$$

$$\frac{\tau_{tr}}{\rho} = \frac{\tau}{\rho} \left[ \frac{h\nu - P_L Y_L \cdot h\bar{\nu}_L}{h\nu} \right]$$